

## Solution Final Exam — Partial Differential Equations

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Duration: 3 hours

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### Question 1 (15 points)

Consider the equation

$$u_x - 3x^2u_y = 0, \quad (1)$$

where  $u = u(x, y)$ .

a. (7 pt) Find the general solution of Eq. (1).

b. (3 pt) Find the solution of Eq. (1) with the auxiliary condition  $u(0, y) = -y^2$ .

Consider now the equation

$$u_x - 3x^2u_y = u. \quad (2)$$

c. (5 pt) Find the general solution of Eq. (2), using the substitution  $u(x, y) = e^x w(x, y)$ .

### Solution

a. We solve the equation for the characteristic curves

$$\frac{dy}{dx} = -3x^2,$$

which directly gives

$$y = -x^3 + C,$$

where  $C$  is the constant of integration. Solving for  $C$  we get

$$C = y + x^3.$$

Since  $y + x^3$  is constant along the characteristic curves we conclude that the solution of the problem has the general form

$$u(x, y) = f(y + x^3),$$

where  $f$  is an arbitrary function of one variable.

b. Applying the general solution we find

$$u(0, y) = f(y) = -y^2.$$

Therefore  $f(s) = -s^2$ , and the solution we are after is

$$u(x, y) = -(y + x^3)^2.$$

c. We have

$$u_x = e^x w + e^x w_x, \quad u_y = e^x w_y.$$

Then Eq. (2) gives

$$e^x(w + w_x - 3x^2 w_y - w) = 0,$$

which can be simplified, to

$$w_x - 3x^2 w_y = 0,$$

which is exactly Eq. (1), and for which we know that the general solution is

$$w = f(y + x^3).$$

Therefore, the general solution for Eq. (2) is

$$u = e^x f(y + x^3),$$

with  $f$  an arbitrary function.

## Question 2 (15 points)

Consider the equation

$$u_{xx} - 4u_{xy} + 4u_{yy} = 0. \quad (3)$$

- a. (3 pt) What is the type (elliptic / hyperbolic / parabolic) of Eq. (3)? Explain your answer.
- b. (8 pt) Find a linear transformation  $(x, y) \rightarrow (s, t)$  that reduces Eq. (3) to one of the standard forms  $u_{ss} + u_{tt} = 0$ ,  $u_{ss} - u_{tt} = 0$ , or  $u_{ss} = 0$ .
- c. (4 pt) Find the general solution of Eq. (3).

### Solution

- a. We have  $a_{11} = 1$ ,  $a_{22} = 4$ , and  $a_{12} = -2$ . Therefore

$$a_{12}^2 = a_{11}a_{22},$$

and Eq. (3) is *parabolic*.

- b. Since the equation is parabolic the standard form is  $u_{ss} = 0$ , or  $\partial_s^2 u = 0$ . Write the original equation as

$$\mathcal{L}u = 0,$$

where

$$\mathcal{L} = \partial_x^2 - 4\partial_x\partial_y + 4\partial_y^2.$$

Then

$$\mathcal{L} = (\partial_x - 2\partial_y)^2,$$

so we can set  $\partial_s = \partial_x - 2\partial_y$  and  $\partial_t = \partial_y$ , that is,

$$\begin{pmatrix} \partial_s \\ \partial_t \end{pmatrix} = \begin{pmatrix} 1 & -2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \partial_x \\ \partial_y \end{pmatrix}.$$

The corresponding coordinate transformation is

$$\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ -2 & 1 \end{pmatrix} \begin{pmatrix} s \\ t \end{pmatrix},$$

or

$$x = s, \quad y = t - 2s,$$

which can be inverted to give

$$s = x, \quad t = y + 2x.$$

- c. We have transformed Eq. (3) to  $u_{ss} = 0$ . The latter has the general solution

$$u(s, t) = f(t) + sg(t).$$

This means that the original equation has the solution

$$u(x, y) = f(y + 2x) + xg(y + 2x).$$

**Question 3 (15 points)**

Consider the equation for the damped string

$$u_{tt} - c^2 u_{xx} + ru_t = 0, \quad (4a)$$

where  $x \in [0, L]$ ,  $t \geq 0$ ,  $c > 0$ ,  $r > 0$ ,  $L > 0$ , and Dirichlet boundary conditions

$$u(0, t) = u(L, t) = 0. \quad (4b)$$

Define the energy of the string as

$$E(t) = \frac{1}{2} \int_0^L (u_t^2 + c^2 u_x^2) dx. \quad (5)$$

- a. (8 pt) Prove that the energy of the string decreases with time, that is,  $dE/dt \leq 0$ .
- b. (7 pt) Prove uniqueness of solutions  $u$  that satisfy Eq. (4a), Eq. (4b), and  $u(x, 0) = g(x)$ ,  $u_t(x, 0) = h(x)$  for  $x \in [0, L]$ .

**Solution**

- a. We compute

$$\frac{dE}{dt} = \frac{1}{2} \int_0^L \frac{\partial}{\partial t} (u_t^2 + c^2 u_x^2) dx = \int_0^L (u_t u_{tt} + c^2 u_x u_{xt}) dx.$$

Using Eq. (4a) we rewrite the last expression as

$$\frac{dE}{dt} = \int_0^L (u_t c^2 u_{xx} + c^2 u_x u_{xt} - r u_t^2) dx = c^2 \int_0^L (u_t u_{xx} + u_x u_{xt}) dx - r \int_0^L u_t^2 dx.$$

Since  $(u_t u_x)_x = u_t u_{xx} + u_x u_{xt}$  we have

$$\frac{dE}{dt} = c^2 \int_0^L (u_t u_x)_x dx - r \int_0^L u_t^2 dx = c^2 [(u_t u_x)|_{x=L} - (u_t u_x)|_{x=0}] - r \int_0^L u_t^2 dx,$$

so

$$\frac{dE}{dt} = c^2 [u_t(L, t) u_x(L, t) - u_t(0, t) u_x(0, t)] - r \int_0^L u_t^2 dx.$$

From the definition of partial derivatives we have for fixed  $x = a$  that

$$u_t(a, t) = \frac{d}{dt} [u(a, t)].$$

In particular,

$$u_t(0, t) = \frac{d}{dt} [u(0, t)] = \frac{d}{dt} [0] = 0,$$

and similarly  $u_t(L, t) = 0$ . Therefore,

$$\frac{dE}{dt} = -r \int_0^L u_t^2 dx.$$

Since  $u_t^2 \geq 0$  we also have  $\int_0^L u_t^2 dx \geq 0$  and since  $r > 0$  we finally get

$$\frac{dE}{dt} \leq 0.$$

- b. Consider two solutions  $u_1, u_2$  that satisfy Eqs. (4a) and (4b). Define  $w = u_1 - u_2$ . Then

$$w_{tt} - c^2 w_{xx} + r w_t = 0,$$

and

$$w(x, 0) = u_1(x, 0) - u_2(x, 0) = g(x) - g(x) = 0,$$

$$w_t(x, 0) = (u_1)_t(x, 0) - (u_2)_t(x, 0) = h(x) - h(x) = 0.$$

Then for the energy  $E(t)$  corresponding to  $w$  we have  $dE/dt \leq 0$  and

$$E(0) = \frac{1}{2} \int_0^L (w_t(x, 0)^2 + c^2 w_x(x, 0)^2) dx.$$

We have  $w_x(x, 0) = \frac{d}{dx}[w(x, 0)] = 0$  and we also saw that  $w_t(x, 0) = 0$ , so

$$E(0) = 0.$$

Since  $dE/dt \leq 0$  we conclude that  $E(t) \leq 0$  for all  $t \geq 0$ . Furthermore, by its definition  $E(t) \geq 0$  so we conclude that  $E(t) = 0$ . This implies that  $w_t(x, t) = w_x(x, t) = 0$  for all  $t \geq 0$  and  $x \in [0, L]$ . Therefore  $w$  is constant and since at  $t = 0$  it is  $w(x, 0) = 0$  we conclude that  $w(x, t) = 0$ . From here  $u_1(x, t) = u_2(x, t)$ .

**Question 4 (20 points)**

Consider the Laplace equation  $\Delta u = 0$  in a half-disk  $H$  of radius  $a$ , that is,

$$H = \{(x, y) : x^2 + y^2 \leq a^2, y \geq 0\},$$

with the boundary conditions  $u(a, \theta) = \sin \theta$  for  $0 \leq \theta \leq \pi$ , and  $u(r, 0) = u(r, \pi) = 0$  for  $0 \leq r \leq a$ .

- a. (5 pt) Separate the Laplace equation in polar coordinates  $r, \theta$  using the ansatz  $u(r, \theta) = R(r)\Theta(\theta)$  and write two ordinary differential equations, one for  $R$  and one for  $\Theta$ .
- b. (7 pt) Solve the eigenvalue equation for  $\Theta$  for the given boundary conditions (find eigenvalues and eigenfunctions). Consider known that the problem has no complex eigenvalues but check for positive, negative, or zero eigenvalues.
- c. (4 pt) Solve the differential equation for  $R$ .
- d. (4 pt) Write the general solution  $u(r, \theta)$  for arbitrary boundary conditions  $u(a, \theta) = h(\theta)$  and then give the solution for the specific boundary conditions in this problem.

**Solution**

- a. Substituting  $u(r, \theta) = R(r)\Theta(\theta)$  into the equation

$$\Delta u = u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta} = 0$$

we get

$$R''\Theta + \frac{1}{r}R'\Theta + \frac{1}{r^2}R\Theta'' = 0.$$

Then

$$r^2\frac{R''}{R} + r\frac{R'}{R} + \frac{\Theta''}{\Theta} = 0.$$

and separating the parts that depend on  $r$  from those that depend on  $\theta$  we find

$$r^2\frac{R''}{R} + r\frac{R'}{R} = -\frac{\Theta''}{\Theta} = \lambda,$$

where  $\lambda$  is a constant. Then we have the two equations

$$r^2R'' + rR' - \lambda R = 0, \quad \Theta'' + \lambda\Theta = 0.$$

- b. The given boundary conditions imply

$$\Theta(0) = \Theta(\pi) = 0.$$

For positive eigenvalues  $\lambda = \beta^2$  we have the solutions

$$\Theta(\theta) = A \cos(\beta\theta) + B \sin(\beta\theta).$$

From here

$$\Theta(0) = A = 0, \quad \Theta(\pi) = A \cos(\beta\pi) + B \sin(\beta\pi) = 0.$$

Then

$$A = 0, \quad B \sin(\beta\pi) = 0,$$

and finally

$$\beta_n = n, n = 1, 2, 3, \dots$$

For  $\lambda = 0$  the solution is

$$\Theta(\theta) = A\theta + B$$

so

$$\Theta(0) = B = 0, \quad \Theta(\pi) = A\pi + B = 0$$

giving the trivial solution  $A = B = 0$  which is rejected.

For  $\lambda = -\gamma^2 < 0$  we have

$$\Theta(\theta) = Ae^{\gamma\theta} + Be^{-\gamma\theta}.$$

Then

$$\Theta(0) = A + B = 0, \quad \Theta(\pi) = Ae^{\gamma\pi} + Be^{-\gamma\pi} = 0,$$

so

$$B = -A, \quad Ae^{-\gamma\pi}(e^{2\gamma\pi} - 1) = 0.$$

The last equation implies that either  $\gamma = 0$  (so  $\lambda = 0$  but we assumed  $\lambda < 0$ ) or  $A = B = 0$  giving the trivial solution so we should also reject the case of negative eigenvalues.

Finally, the eigenvalues are

$$\lambda_n = \beta_n^2 = n^2,$$

and the eigenfunctions

$$\Theta_n(\theta) = \sin(n\theta).$$

- c. The differential equation for  $R$  is  $r^2 R'' + rR' - n^2 R = 0$ . Try the solution  $R = r^\alpha$  which gives  $\alpha = \pm n$ . Therefore

$$R_n(r) = A_n r^n + B_n r^{-n}.$$

Since we want the solutions to be well-defined at  $r = 0$  we set  $B_n = 0$ .

- d. The general solution is

$$u(r, \theta) = \sum_{n=1}^{\infty} A_n r^n \sin(n\theta).$$

For  $r = a$  we have

$$u(a, \theta) = \sum_{n=1}^{\infty} A_n a^n \sin(n\theta) = h(\theta)$$

so

$$A_n = \frac{2}{\pi a^n} \int_0^\pi h(\theta) \sin(n\theta).$$

In the specific case here we have

$$h(\theta) = \sin \theta = \sum_{n=1}^{\infty} A_n a^n \sin(n\theta),$$

and comparing the two expressions we see that  $A_1 = 1/a$ ,  $A_n = 0$  in all other cases. Therefore the solution for the specific boundary conditions is

$$u(r, \theta) = \frac{r}{a} \sin \theta.$$

**Question 5 (10 points)**

Suppose that  $u$  is a harmonic function in the closed disk  $D = \{r \leq 1\}$  and that  $u = 2 \cos(3\theta) + 1$  for  $r = 1$ .

- a. (5 pt) What are the maximum and minimum values of  $u$  in  $D$ ?
- b. (5 pt) Find the value of  $u$  at the origin.

**Solution**

- a. Since  $u$  is harmonic it attains its maximum and minimum values at the boundary. At the boundary we have  $u = 2 \cos(3\theta) + 1$ . Since  $-1 \leq \cos(3\theta) \leq 1$  and the values  $\pm 1$  are attained for  $\theta \in [0, 2\pi]$  we conclude that  $-1 \leq u \leq 3$  at the boundary and  $u$  attains the maximum value 3 and the minimum value  $-1$  at some point on the boundary. Therefore these are also the respective maximum and minimum values on  $D$ .
- b. Poisson's formula is

$$u(r, \theta) = \frac{a^2 - r^2}{2\pi} \int_0^{2\pi} \frac{h(\phi)}{a^2 - 2ar \cos(\theta - \phi) + r^2} d\phi.$$

Applying Poisson's formula for  $r = 0$ ,  $a = 1$ ,  $h(\phi) = 2 \cos(3\phi) + 1$  gives

$$u(0) = \frac{1}{2\pi} \int_0^{2\pi} (2 \cos(3\phi) + 1) d\phi = 1.$$



**Question 6 (15 points)**

Consider the function

$$f(x) = \pi - x, \quad \text{with } x \in [0, \pi],$$

and its Fourier sine series.

- a. (2 pt) Does the Fourier sine series converge in the  $L^2$  sense? Explain your answer.
- b. (3 pt) What is the pointwise limit of the Fourier sine series for  $x \in \mathbb{R}$ ?
- c. (3 pt) How does the Gibbs phenomenon manifest itself in the Fourier sine series? That is, at which point(s) in  $[0, \pi]$  the Gibbs phenomenon appears and approximately how much is the “overshoot” there?
- d. (7 pt) Compute the coefficients of the Fourier sine series for  $f(x)$ .

**Solution**

- a. The function  $f$  is bounded in  $[0, \pi]$ , therefore

$$\|f\|^2 = \int_0^\pi f(x)^2 dx < +\infty.$$

This means that the Fourier sine series converges in the  $L^2$  sense.

- b. The pointwise limit of the Fourier sine series can be deduced from the odd-periodic extension  $f_{\text{ext}}(x)$  of  $f(x)$  from  $[0, \pi]$  to  $\mathbb{R}$ . This is constructed by first considering the extension of  $f$  to an odd function defined in  $[-\pi, \pi]$  and then the further periodic extension to  $\mathbb{R}$ . This extension  $f_{\text{ext}}(x)$  is discontinuous at  $x = 2k\pi$ ,  $k \in \mathbb{Z}$  and  $f_{\text{ext}}(2k\pi^+) = \pi$  while  $f_{\text{ext}}(2k\pi^-) = -\pi$ . Therefore the Fourier sine series converges pointwise at  $x = 2k\pi$  to  $\frac{1}{2}[f_{\text{ext}}(2k\pi^+) + f_{\text{ext}}(2k\pi^-)] = 0$ . At all other  $x \in \mathbb{R}$ ,  $f_{\text{ext}}(x)$  is continuous so the Fourier series converges pointwise to  $f_{\text{ext}}(x)$ .
- c. The odd-periodic extension  $f_{\text{ext}}$  of  $f$  is discontinuous at  $x = 2k\pi$ ,  $k \in \mathbb{Z}$  so the only point in  $[0, \pi]$  where  $f_{\text{ext}}$  is discontinuous is  $x = 0$ . The jump of  $f_{\text{ext}}$  at  $x = 0$  is  $f_{\text{odd}}(0^+) - f_{\text{odd}}(0^-) = 2\pi$ . Therefore, for  $x \in [0, \pi]$  the Gibbs phenomenon appears at  $x = 0$  and the overshoot is approximately  $0.09 \cdot (2\pi) \simeq 0.56$ .
- d. We have

$$A_n = \frac{2}{\pi} \int_0^\pi (\pi - x) \sin(nx) dx = \frac{2}{\pi} \int_0^\pi (\pi - x) \left( -\frac{1}{n} \cos(nx) \right)' dx.$$

Integration by parts gives

$$A_n = -\frac{2}{n\pi} \left[ (\pi - x) \cos(nx) \right]_0^\pi - \frac{2}{n\pi} \int_0^\pi \cos(nx) dx = -\frac{2}{n\pi} \left[ (\pi - x) \cos(nx) + \frac{\sin(nx)}{n} \right]_0^\pi.$$

Then

$$A_n = \frac{2}{n}.$$